

Chaotic capture of (dark) matter by binary systems

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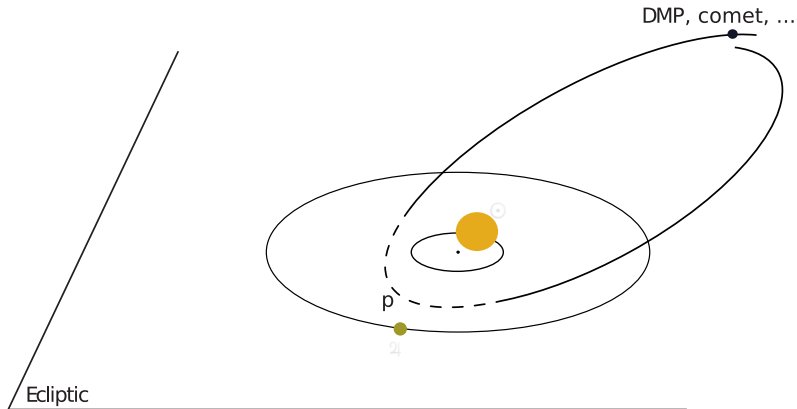
– Journées du GRAM, Besançon 2018 –

References :

- J. L., D. Shepelyansky, **Dark matter chaos in the Solar system**, MNRAS Letters 430, L25-L29 (2013)
G. Rollin, J. L., D. Shepelyansky, **Chaotic enhancement of dark matter density in binary systems**, A&A 576, A40 (2015)
P. Haag, G. Rollin, J. L., **Symplectic map description of Halley's comet dynamics**, Physics Letters A 379 (2015) 1017-1022
J. L., D. Shepelyansky, I. Shevchenko, **Kepler map**, Scholarpedia, 13(2) :33238 (2018)



(Dark) matter capture – Three-body problem



Possible DMP capture (or comet capture) due to Jupiter and Sun rotations around the SS barycenter.

Dark matter capture – Restricted circular three-body problem

$$m_{\text{DMP}} \ll m_+ \ll m_\odot$$

Dark matter capture – Restricted circular three-body problem

Newton's equations

$$m_{\text{DMP}} \ll m_{\text{p}} \ll m_{\odot}$$

$$\ddot{\mathbf{r}} = \frac{1 - m_{\text{p}}}{\|\mathbf{r}_{\odot}(t) - \mathbf{r}\|^3} (\mathbf{r}_{\odot}(t) - \mathbf{r}) + \frac{m_{\text{p}}}{\|\mathbf{r}_{\text{p}}(t) - \mathbf{r}\|^3} (\mathbf{r}_{\text{p}}(t) - \mathbf{r})$$

$$G = 1, \quad m_{\text{p}} + m_{\odot} = 1, \quad \|\dot{\mathbf{r}}_{\text{p}}\| \simeq 13\text{km.s}^{-1} = 1, \quad \|\mathbf{r}_{\text{p}}\| = 1$$

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Energy change after a passage at perihelion (wide encounter)

$$F \sim \frac{m_{\text{p}}}{m_{\odot}} \|\dot{\mathbf{r}}_{\text{p}}\|^2 \simeq 10^{-3}$$

Dark matter capture – Restricted circular three-body problem

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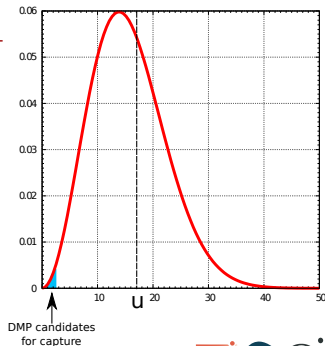
Energy change after a passage at perihelion (wide encounter)

$$F \sim \frac{m_{\gamma}}{m_{\odot}} \|\dot{\mathbf{r}}_{\gamma}\|^2 \simeq 10^{-3}$$

Assuming a Maxwellian distribution of Galactic DMP velocities

$$f(v)dv \sim v^2 \exp\left(-3v^2/2u^2\right) dv$$

with $u \simeq 220 \text{ km.s}^{-1} \sim 17$ (mean DMP velocity)



Dark matter capture – Restricted circular three-body problem

Newton's equations

$$m_{\text{DMP}} \ll m_{\text{d}} \ll m_{\odot}$$

$$\ddot{\mathbf{r}} = \frac{1 - m_{\text{d}}}{\|\mathbf{r}_{\odot}(t) - \mathbf{r}\|^3} (\mathbf{r}_{\odot}(t) - \mathbf{r}) + \frac{m_{\text{d}}}{\|\mathbf{r}_{\text{d}}(t) - \mathbf{r}\|^3} (\mathbf{r}_{\text{d}}(t) - \mathbf{r})$$

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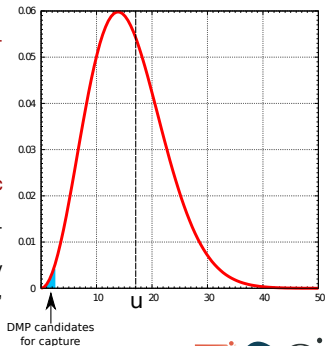
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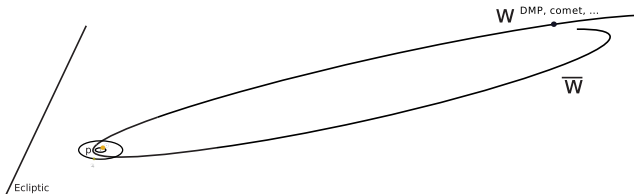
As $F \ll u^2$, not many candidates for capture among Galactic DMPs

Most of the capturable DMPs have close to parabolic approaching trajectories ($E \sim 0$)

Direct simulation of Newton's equations is difficult : very elongated ellipses, not many particles can be simulated, CPU time consuming (Peter 2009)



Kepler map



x : Jupiter's phase when particle at pericenter ($x = \varphi/2\pi \mod 1$)

w : particle energy at apocenter ($w = -2E/m_{\text{DMP}}$)

Symplectic Kepler map

$$\begin{aligned}\bar{w} &= w + F(x) &= w + W \sin(2\pi x) &\leftarrow \text{energy change after a kick} \\ \bar{x} &= x + \bar{w}^{-3/2} &&\leftarrow \text{third Kepler's law}\end{aligned}$$

Map already used in the study of :

- ▶ Cometary clouds in Solar systems (Petrosky 1986)
- ▶ Chaotic dynamics of Halley's comet (Chirikov & Vecheslavov 1986)
- ▶ Microwave ionization of hydrogen atoms (see e.g. Shepelyansky, scholarpedia)

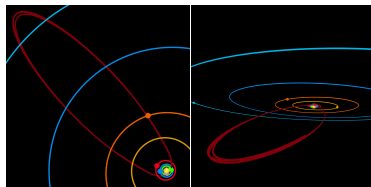
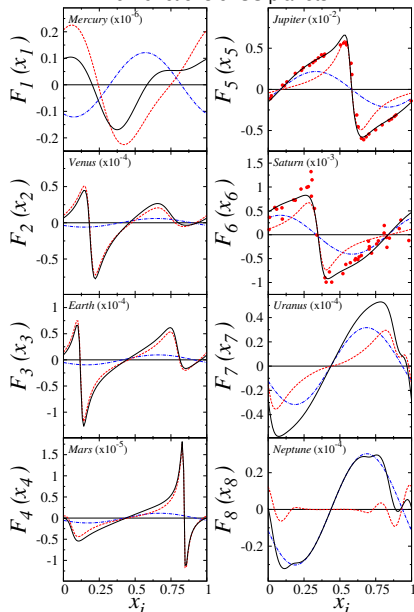
Advantage : if the kick function $F(x)$ is known the dynamics of a huge number of particles can be simulated.



Let's make a digression ...

Halley map – Cometary case

Kick functions of SS planets



Rollin, Haag, J. L., Phys. Lett. A 379 (2015) 1017-1022

Our calculations match direct observation data and previous numerical data (Yeomans & Kiang 1981)

$$F(x_1, \dots, x_8) \simeq \sum_{i=1}^8 F_i(x_i)$$

$$F_i(x_i) = -2\mu_i \int_{-\infty}^{+\infty} \nabla \left(\frac{\mathbf{r} \cdot \mathbf{r}_i}{r^3} - \frac{1}{\|\mathbf{r} - \mathbf{r}_i\|} \right) \cdot \dot{\mathbf{r}} dt$$

Two main contributions

- Direct planetary Keplerian potential
- Rotating gravitational dipole potential due to the Sun movement around Solar System barycenter

Halley map – Cometary case

Renormalized kick function

$$f_i(x_i) = F_i(x_i)/v_i^2/\mu_i$$

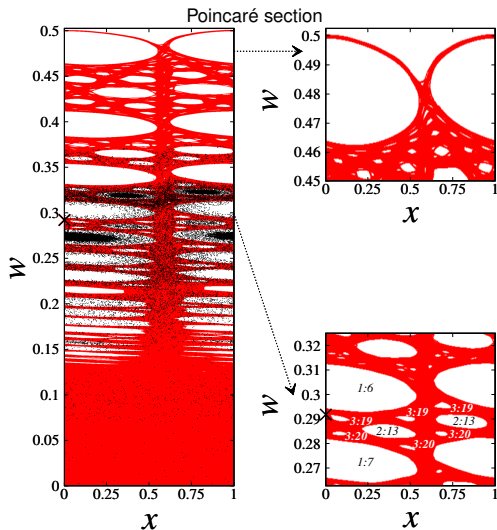
Exponential decay with q for $q > 1.5a_i$
More precisely

$$f_i \simeq 2^{1/4} \pi^{1/2} \left(\frac{q}{a_i} \right)^{-1/4} \exp \left(-\frac{2^{3/2}}{3} \left(\frac{q}{a_i} \right)^{3/2} \right)$$

(Heggie 1975, Petrosky 1986, Petrosky & Broucke 1988, Roy & Haddow 2003, Shevchenko 2011, J.L., Shepelyansky, Shevchenko 2017)



Halley map – Dynamical chaos



Symplectic Halley map

$$\begin{aligned}\bar{x} &= x + \bar{w}^{-3/2} \\ \bar{w} &= w + F(x)\end{aligned}$$

Chaotic dynamics of 1P/Halley

“Lifetime” $\sim 10^7$ years

Let's come back to (dark) matter ...

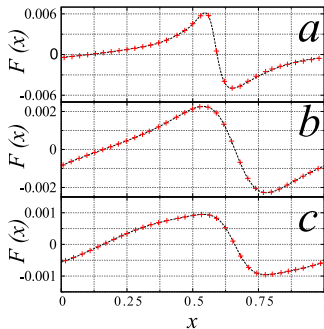
Dark matter capture – Dark map

Kick function determination

We determine numerically the kick function for any parabolic orbit (q, i, ω)

$$F(x) = F_{q,i,\omega}(x)$$

By nonlinear fit we obtain analytical functions.



a : Halley's comet

b : $q = 1.5, \omega = 0.7, i = 0$

c : $q = 0.5, \omega = 0., i = \pi/2$

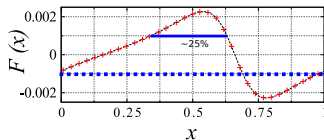
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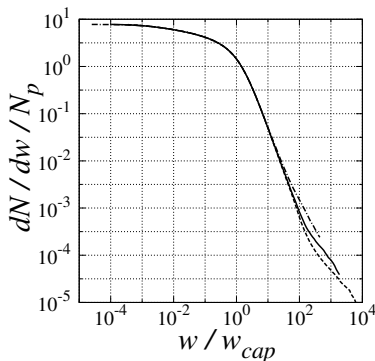
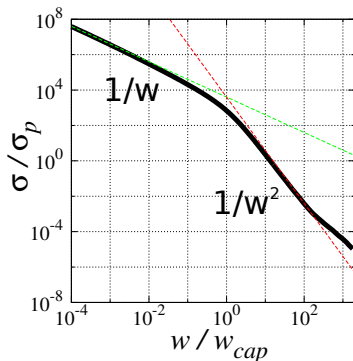
Chance to be captured with a given energy ($w < 0 \leftrightarrow E > 0$)

$$h_{q,i,\omega}(w)$$

Dark matter – Capture cross section

$$w_{cap} = \frac{m_{\gamma}}{m_{\odot}} \|\dot{\mathbf{r}}_{\gamma}\|^2 \simeq 10^{-3}$$

$$\sigma_p = \pi \|\mathbf{r}_{\gamma}\|^2 \quad \text{area enclosed by Jupiter's orbit}$$



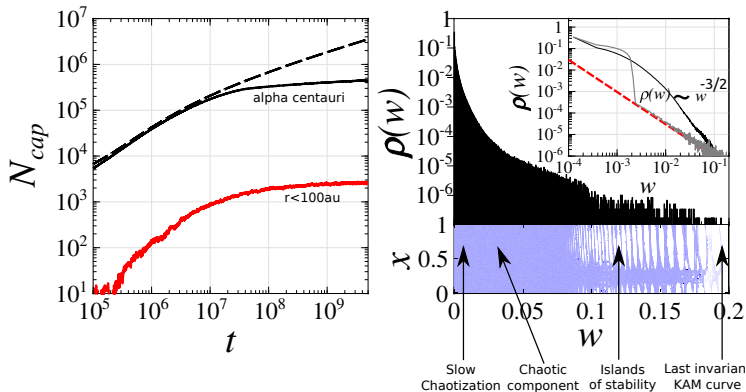
$$\sigma/\sigma_p \simeq \pi \frac{m_{\odot}}{m_{\gamma}} \frac{w_{cap}}{|w|} \quad \text{in agreement with Khriplovich \& Shepelyansky 2009}$$

- Predominance of wide encounters as suggested by Peter 2009
- Very small contribution from close encounters invalidating previous numerical results (Gould & Alam 2001 and Lundberg & Edsjö 2004)

Dark map – Dark matter capture

Simulation of the (isotropic) injection, the capture and the escape of DMPs during the whole lifetime of the Solar system.

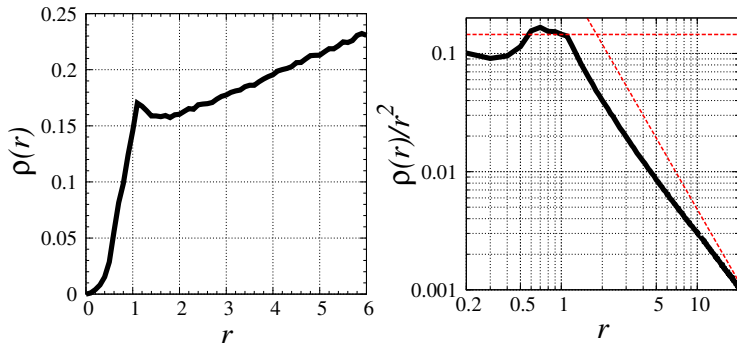
Injection of $N_{tot} \simeq 1.5 \times 10^{14}$ DMPs with energy $|w|$ in the range $[0, \infty]$ with $N_H = 4 \times 10^9$ DMPs in the Halley's comet energy interval $[0, w_H]$.



- Equilibrium reached after a time $t_d \sim 10^7 \text{ yr}$ similar to the diffusive escape time scale of the Halley's comet (Chirikov & Vecheslavov 1989) → Equilibrium energy distribution $\rho(w)$
- The dynamics of dark matter particles in the Solar system is essentially chaotic

Back to real space – Density distribution of captured DMPs

Nowadays equilibrium density distribution ($t_S = 4.5 \times 10^9 \text{ yr}$)



- **The profile of the radial density $\rho(r) \propto dN/dr$ is similar to those observed for galaxies where DMP mass is dominant.** Indeed $\rho(r)$ is almost flat (increases slowly) right after Jupiter orbit ($r = 1$) $\rightarrow$ according to virial theorem the circular velocity of visible matter is consequently constant as observed e.g. in Rubin 1980

Virial theorem : $v_m^2 \sim \int_0^r dr' \rho(r')/r \sim \rho(r) \sim r^2 (\rho(r)/r^2) \underset{\text{here}}{\sim} r^2 r^{-1.53} \sim r^{1/2}$

Ergodicity along radial dynamics : $d\mu \sim dN \sim \rho(r) dr \sim dt \sim dr/v_r \sim r^{1/2} dr$

Consequently, $v_m \propto r^{0.25}$ (Dark map) to compare to $v_m \propto r^{0.35}$ (Rubin 1980)

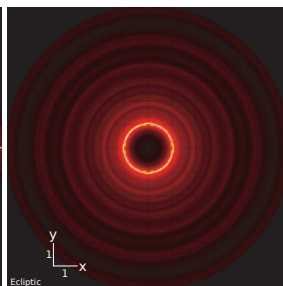
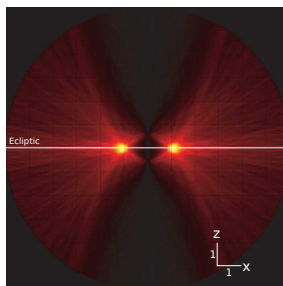
Back to real space – Density distribution of captured DMPs

Surface density

$$\rho_s(z, R) \propto dN/dz dR$$

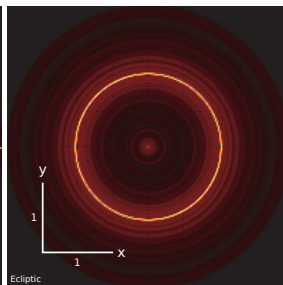
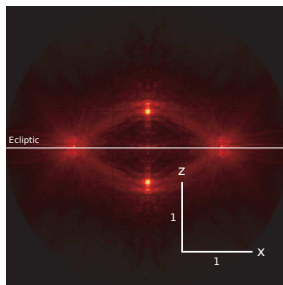
where

$$R = \sqrt{x^2 + y^2}$$



Volume density

$$\rho_v(x, y, z) \propto dN/dxdydz$$



How much dark matter is present in the Solar system ?

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The total mass of DMP passed through the System solar during its lifetime $t_S = 4.5 \times 10^9 \text{yr}$ is

$$M_{\text{tot}} = \rho_g t_S \int_0^\infty dv v f(v) \sigma(v) \approx 35 \rho_g t_S G \|\mathbf{r}_{\gamma_+}\| M_\odot / u \approx 0.9 \times 10^{-6} M_\odot \sim M_\oplus$$

At time t_S the mass of captured DMPs in the Solar system is

$$\begin{aligned} M_{AC} &\approx \eta_{AC} M_{\text{tot}} \approx 2 \times 10^{-15} M_\odot && \text{within } r < 0.5 \text{ distance}_{\text{Sun}-\alpha \text{Centauri}} \\ M_{100\text{au}} &\approx \eta_{100\text{au}} M_{\text{tot}} \approx 1.3 \times 10^{-17} M_\odot && \text{within } r < 100\text{au} \end{aligned}$$

The captured DMP mass in the volume of the Neptune orbit radius is

$$M_\Psi \approx \eta_\Psi M_{AC} \approx 0.9 \times 10^{-18} M_\odot \approx 1.5 \times 10^{15} \text{g}$$

The captured DMP mass in the volume of the Jupiter orbit radius is

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The average volume density of captured dark matter inside the Jupiter orbit sphere is

$$\rho_{\gamma_+} = \frac{3M_{\gamma_+}}{4\pi r_{\gamma_+}^3} \approx 5 \times 10^{-29} \text{g/cm}^3 \approx 1.2 \times 10^{-4} \rho_g \ll \rho_g \text{ (Galactic DMP density)}$$

Globally, not much dark matter captured by the Solar system, but ...

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Let's compare to the capturable DMP density

$$\rho_{gH} = \rho_g \int_0^{\sqrt{wH}} dv v f(v) \approx 1.4 \times 10^{-32} \text{g/cm}^3 \quad \Rightarrow$$

Huge chaotic enhancement $\zeta = \rho_{\gamma_+} / \rho_{gH} \approx 4 \times 10^3$ of the density of actually capturable DMPs.

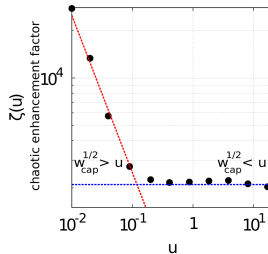
The long range interaction capture mechanism is very efficient for binary systems (1+2) with

$$m_1 \gg m_2$$



(Dark) matter capture in binary systems

$$w_{\text{cap}} \sim w_H \sim 0.005$$
$$u = 220 \text{ km/s}$$

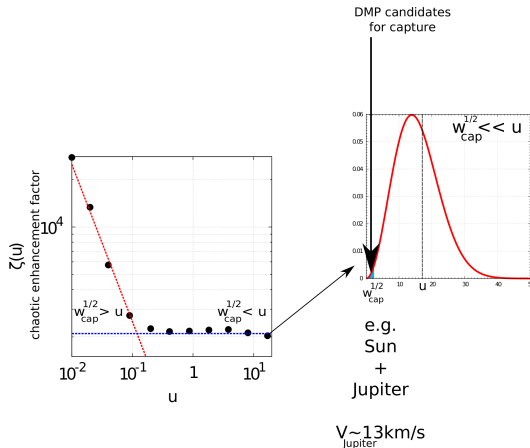


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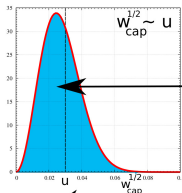


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(Dark) matter capture in binary systems

e.g.
Black hole
+
Star companion

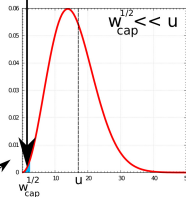
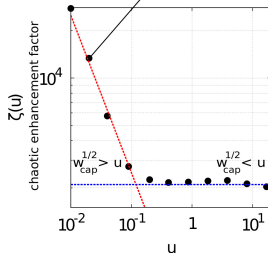
$$V_{\text{star}} \sim c/40$$



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DMP candidates
for capture



e.g.
Sun
+
Jupiter

$$V_{\text{Jupiter}} \sim 13 \text{ km/s}$$

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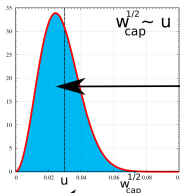
(Dark) matter capture in binary systems

Global
volume density
enhancement
 $\times 10^4$

**all the galactic
DMPs are captured**

e.g.
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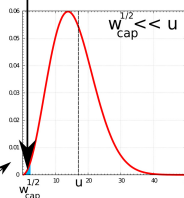
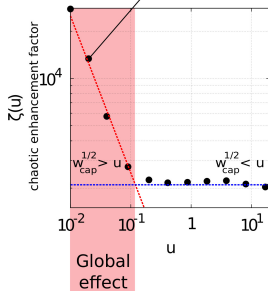
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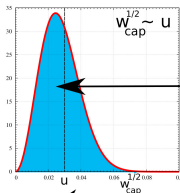
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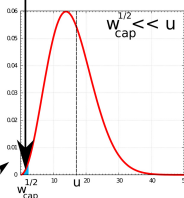
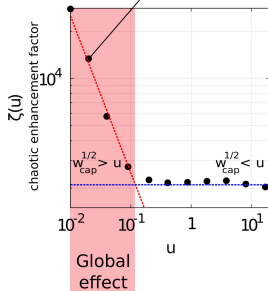
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Ionization process:
- DMPs acceleration
- wandering black holes
- ...

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Thank You !