

# Chaotic dark matter in the Solar system and galaxies

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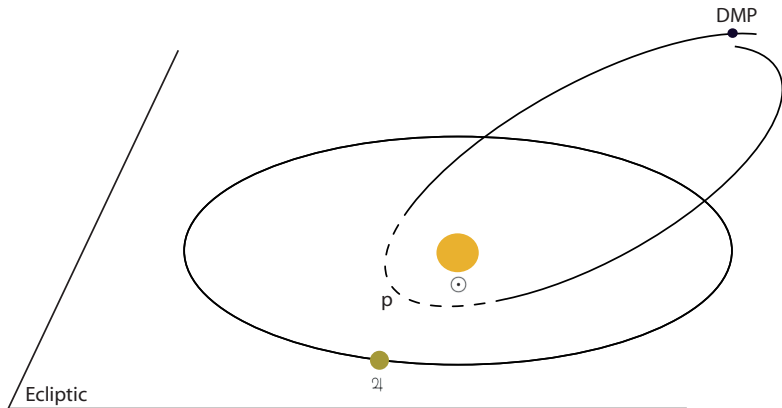
## References :

J. L., D. Shepelyansky, **Dark matter chaos in the Solar system**, MNRASL 430, L25-L29 (2013), arXiv :1211.0903

G. Rollin, J. L., D. Shepelyansky, **Chaotic enhancement of dark matter density in binary systems and galaxies**, to be submitted



# Dark matter capture – Three-body problem



Possible DMP capture due to Jupiter rotation around the Sun

# Dark matter capture – Restricted circular three-body problem

$$m_{\text{DMP}} \ll m_{\oplus} \ll m_{\odot}$$

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Newton's equations

$$m_{\text{DMP}} \ll m_{\gamma} \ll m_{\odot}$$

$$\ddot{\mathbf{r}} = \frac{1 - m_{\gamma}}{\|\mathbf{r}_{\odot}(t) - \mathbf{r}\|^3} (\mathbf{r}_{\odot}(t) - \mathbf{r}) + \frac{m_{\gamma}}{\|\mathbf{r}_{\gamma}(t) - \mathbf{r}\|^3} (\mathbf{r}_{\gamma}(t) - \mathbf{r})$$

$$G = 1, \quad m_{\gamma} + m_{\odot} = 1, \quad \|\dot{\mathbf{r}}_{\gamma}\| \simeq 13\text{km}\cdot\text{s}^{-1} = 1, \quad \|\mathbf{r}_{\gamma}\| = 1$$

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Energy change after a passage at perihelion (in absence of close encounter)

$$F \sim \frac{m_{\text{p}}}{m_{\odot}} \|\dot{\mathbf{r}}_{\text{p}}\|^2 \simeq 10^{-3} \quad (\text{Petrosky 86'})$$

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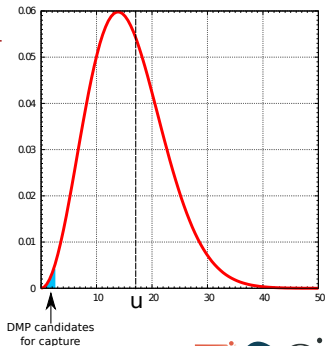
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Assuming a Maxwellian distribution of Galactic DMP velocities

$$f(v)dv \sim v^2 \exp\left(-3v^2/2u^2\right) dv$$

with  $u \simeq 220\text{km}\cdot\text{s}^{-1} \sim 17$  (mean DMP velocity)



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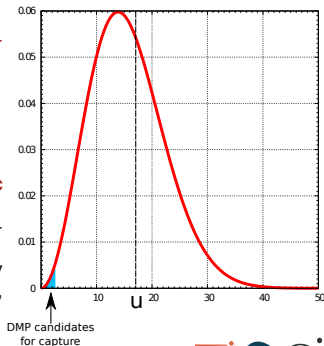
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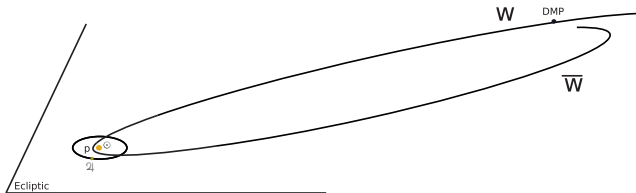
As  $F \ll u^2$ , not many candidates for capture among Galactic DMPs

Most of the capturable DMPs have close to parabolic approaching trajectories ( $E \sim 0$ )

Direct simulation of Newton's equations is difficult : very elongated ellipses, not many particles can be simulated, CPU time consuming (Peter 09')



# Dark matter capture – Kicked model

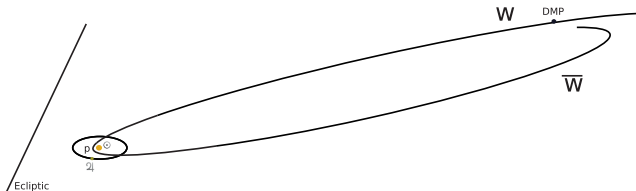


$x$  : Jupiter's phase when DMP at perihelion ( $x = \varphi/2\pi \pmod{1}$ )

$w$  : DMP energy ( $w = -2E/m_{\text{DMP}}$ )



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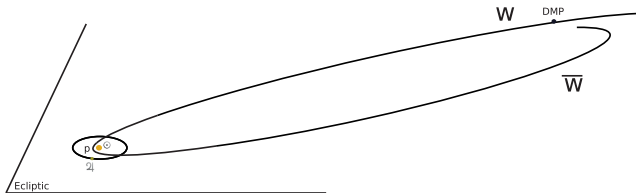
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## Symplectic Kepler-Petrosky map

$$\begin{aligned}\bar{x} &= x + \bar{w}^{-3/2} && \text{third Kepler's law} \\ \bar{w} &= w + F(x) && \text{energy change after a kick}\end{aligned}$$

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Map already used in the study of :

- ▶ Cometary clouds in Solar systems (Petrosky 86')
- ▶ Chaotic dynamics of Halley's comet (Chirikov, Vecheslavov 89')
- ▶ Microwave ionization of hydrogen atoms (see e.g. Shepelyansky, scholarpedia)

Advantage : providing the fact the kick function  $F(x)$  is known the dynamics of a huge number of particles can be simulated.



# Dark matter capture – Dark map

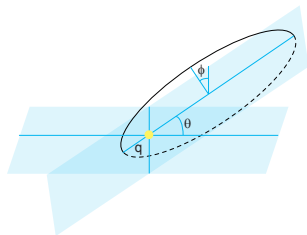
## Kick function determination

$$F(x) \rightarrow F_{\ell, \theta, \phi}(x) \sim F_{q, \theta, \phi}(x)$$

The kick function is different for each approaching configuration  $(\ell, \theta, \phi) \sim (q, \theta, \phi)$

For close to 1 eccentricities the perihelion is such as

$$q \sim \frac{\ell^2}{2}$$



# Dark matter capture – Dark map

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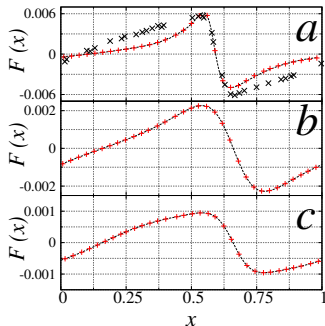
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For each approaching configuration  $(q, \theta, \phi)$ , the integration of the three-body problem gives the kick function.



a : Halley's comet

b :  $q = 1.5, n = 4, \theta = 0.7, \phi = 0$

c :  $q = 0.5, n = 4, \theta = 0., \phi = \pi/2$

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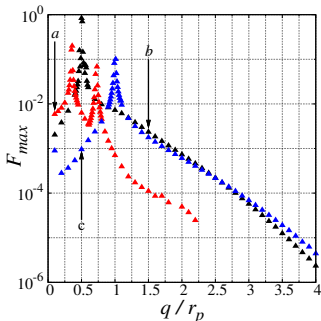
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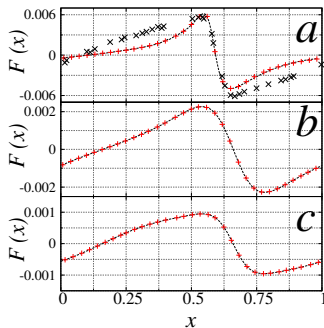
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Exponential decay in agreement with Petrosky 86'



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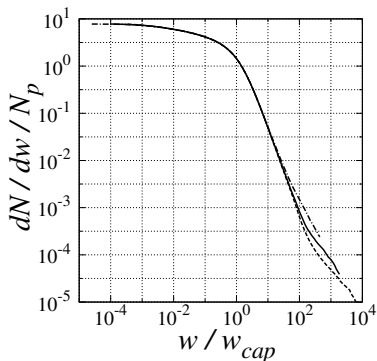
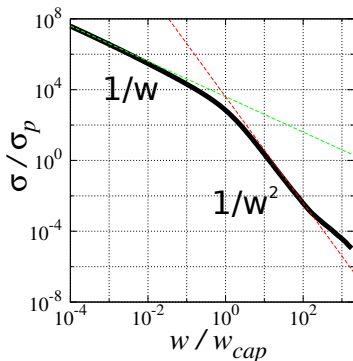
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# Dark matter – Capture cross section

$$w_{cap} = \frac{m_{\eta_+}}{m_{\odot}} \|\dot{\mathbf{r}}_{\eta_+}\|^2 \simeq 10^{-3}$$

$$\sigma_p = \pi \|\mathbf{r}_{\eta_+}\|^2 \quad \text{Jupiter orbit area}$$



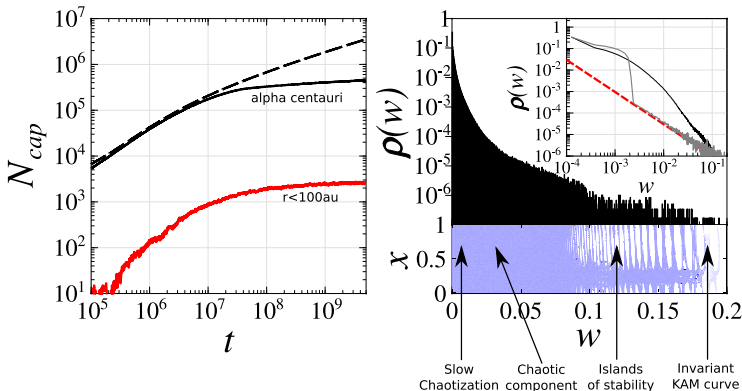
$$\sigma/\sigma_p \simeq \pi \frac{m_{\odot}}{m_{\eta_+}} \frac{w_{cap}}{w} \quad \text{in agreement with Khriplovich \& Shepelyansky 09'}$$

- ▶ Predominance of long range interaction as suggested by Peter 09'
- ▶ Very small contribution from close encounters invalidating previous numerical results (Gould & Alam 01' and Lundberg & Edsjö 04')

# Dark matter evolution – Chaotic dynamics

Simulation of the (isotropic) injection, the capture and the escape of DMPs during the whole lifetime of the Solar system.

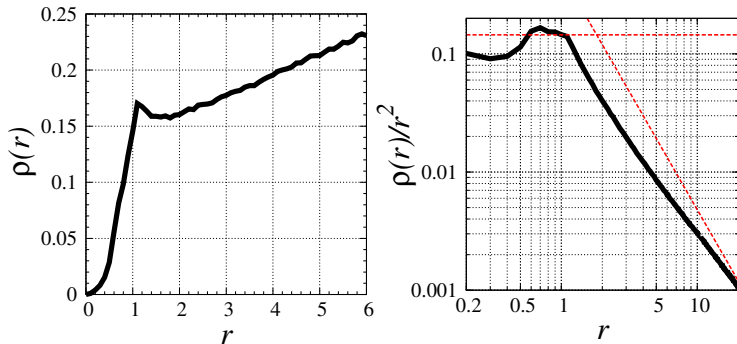
Injection of  $N_{tot} \simeq 1.5 \times 10^{14}$  DMPs with energy  $|w|$  in the range  $[0, \infty]$  with  $N_H = 4 \times 10^9$  DMPs in the Halley's comet energy interval  $[0, w_H]$ .



- ▶ Equilibrium reached after a time  $t_d \sim 10^7$  yr similar to the diffusive escape time scale of the Halley's comet (Chirikov & Vecheslavov 89')  $\rightarrow$  Equilibrium energy distribution  $\rho(w)$
- ▶ The dynamics of dark matter particles in the Solar system is essentially chaotic

# Back to real space – Density distribution of captured DMPs

Nowadays equilibrium density distribution ( $t_S = 4.5 \times 10^9 \text{ yr}$ )



- **The profile of the radial density  $\rho(r) \propto dN/dr$  is similar to those observed for galaxies where DMP mass is dominant.** Indeed  $\rho(r)$  is almost flat (increases slowly) right after Jupiter orbit ( $r = 1$ )  $\rightarrow$  according to virial theorem the circular velocity of visible matter is consequently constant as observed e.g. in Rubin 80'  
More precisely,  $v_m \propto r^{0.25}$  (Dark map) quite close to  $v_m \propto r^{0.35}$  (Rubin 80')



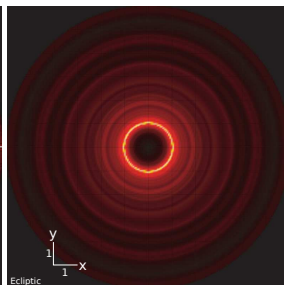
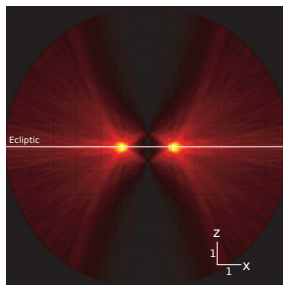
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## Surface density

$$\rho_s(z, R) \propto dN/dz dR$$

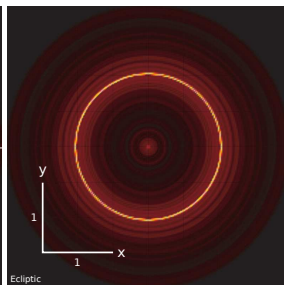
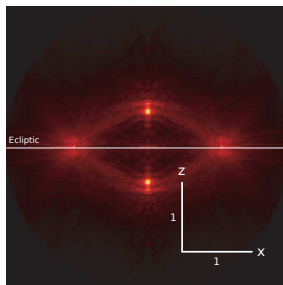
where

$$R = \sqrt{x^2 + y^2}$$



## Volume density

$$\rho_v(x, y, z) \propto dN/dxdydz$$



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The total mass of DMP passed through the System solar during its lifetime  $t_S = 4.5 \times 10^9$ yr is

$$M_{\text{tot}} = \rho_g t_S \int_0^\infty dv v f(v) \sigma(v) \approx 35 \rho_g t_S G \|\mathbf{r}_\oplus\| M_\odot / u \approx 0.9 \times 10^{-6} M_\odot \sim M_\oplus$$

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$$\begin{aligned} M_{AC} &\approx \eta_{AC} M_{\text{tot}} \approx 2 \times 10^{-15} M_\odot && \text{within } r < 0.5 \text{ distance}_{\text{Sun}-\alpha\text{Centauri}} \\ M_{100\text{au}} &\approx \eta_{100\text{au}} M_{\text{tot}} \approx 1.3 \times 10^{-17} M_\odot && \text{within } r < 100\text{au} \end{aligned}$$

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The captured DMP mass in the volume of the Neptune orbit radius is

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Let's compare to the capturable DMP density

$$\rho_{gH} = \rho_g \int_0^{\sqrt{wH}} dv v f(v) \approx 1.4 \times 10^{-32} \text{g/cm}^3$$

**Huge chaotic enhancement  $\zeta = \rho_{\gamma_+} / \rho_{gH} \approx 4 \times 10^3$  of the density of actually capturable DMPs.**



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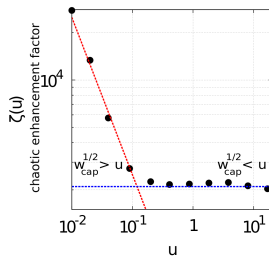
**The long range interaction capture mechanism is very efficient for binary systems (1+2) with**

$$m_1 \gg m_2$$



# Conclusion : Dark matter capture in binary systems (preliminary results)

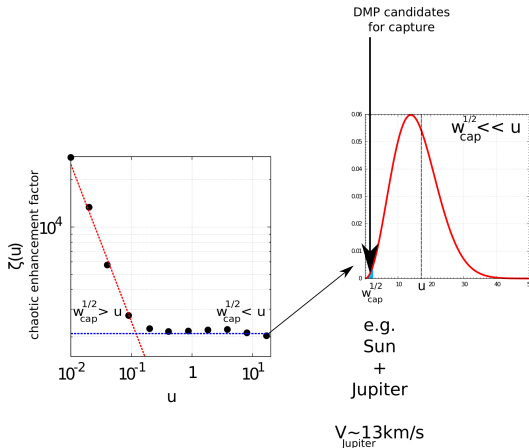
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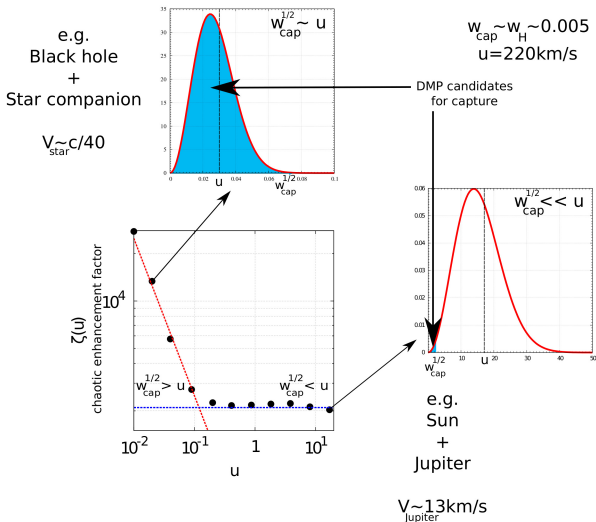
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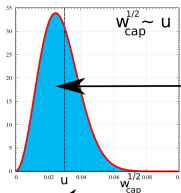
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Global volume density enhancement  $\times 10^4$

**all the galactic DMPs are captured**

e.g.  
Black hole  
+  
Star companion

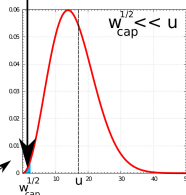
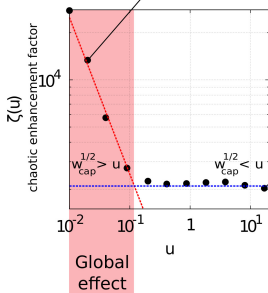
$$V_{\text{star}} \sim c/40$$



$$w_{\text{cap}} \sim w_H \sim 0.005$$

$$u = 220 \text{ km/s}$$

DMP candidates for capture



e.g.  
Sun  
+  
Jupiter

$$V_{\text{Jupiter}} \sim 13 \text{ km/s}$$

Volume density enhancement of capturable DMPs  $\times 10^3$

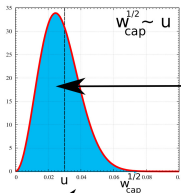
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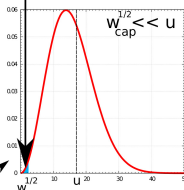
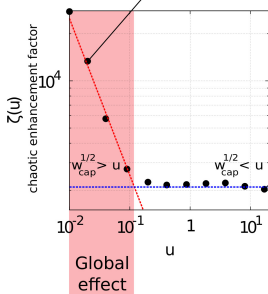
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Volume density enhancement of capturable DMPs  $\times 10^3$

Ionization process:  
- DMPs acceleration  
- wandering black holes  
- ...

# Thank You !